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Study of the Static Response of a Coated Microbubble Under the AFM: Numerical & Asymptotic Analysis

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Outline

- Introduction
- Mathematical Modeling
- Numerical Analysis
- Results
- Conclusions

- Elastic coating (Polymer or Lipid)
- Initial Diameter ~ 2-5 μm
- Initial Thickness ~ 1-30 nm
- Strong Backscattered Acoustic Signal
- Biocompatibility
- Medical Imaging-Diagnostic
 Applications (Heart, Kidney, Liver)





[Blomley et al., Br. Med. J., 2001]

Treatment with microbubbles Drug or gene delivery



[Ferrara et al., Annu. Rev. Biomed. Eng., 2007]





Experimental Investigation of the Mechanical Properties with the Atomic Force Microscope (AFM).

Microbubble covered with polymer





Scope of the Present Work

- Simulation of the static response of CA microbubbles.
- Comparison with AFM experimental measurements.
- Estimation of the elastic properties (Young's & bending modulus).
- Optimization of the design of microbubbles

Classic Contact Problem (Polymeric Bubbles)



Comparison against experiments

[Pozrikidis, J. Fluid Mech., 2001]
[Tsiglifis & Pelekasis, Phys. Fluids, 2011]
[Barthès-Biesel et al., J. Fluid Mech., 2002]
[Updike & Kallins, Trans. ASME, 1970, 1972, 1972]

Normal and Tangential Force Balance $\Delta \vec{F} = \Delta \vec{P} \Rightarrow \begin{cases} \Delta \vec{F}_{N} = (P_{G} - P_{atm})I \cdot \vec{n} \\ \Delta \vec{F}_{s} = \vec{0} \end{cases}$ $\Delta \vec{F} = -\vec{\nabla}_{s} \cdot (\tau_{ss}\vec{t}_{s}\vec{t}_{s} + \tau_{\varphi\varphi}\vec{t}_{\varphi}\vec{t}_{\varphi} + q\vec{t}_{s}\vec{n}) + \sigma\vec{\nabla}_{s} \cdot \vec{n}$ $\vec{q} = \vec{\nabla}_{s} \cdot \underline{m} \cdot (\underline{I} - \vec{n}\vec{n}) \quad ; \quad q^{+}(\xi_{c}) + q^{-}(\xi_{c}) = P_{ext}$ Kinematic Condition $z_{\xi} = 0 \quad \text{at} \quad \xi = \xi_{c}$ Isothermal Gas Compression

 $P_{G0}V_0^{\gamma} = P_GV^{\gamma}$

Constitutive Equation: Hook's Law $\tau = G_s \frac{1+\nu}{1-\nu} (\lambda^2 - 1)$ $m_i = k_b (K_i + \nu K_j) / \lambda_j, \quad K_i \equiv \lambda_i k_i - k_i^R$

Effect of pre-stress $\lambda_i = \frac{dS}{dS^{SF}}, \text{ where } \lambda_i^0 = \frac{dS^0}{dS^{SF}} \neq 1$ Boundary Conditions

$$r_{\xi} = \vartheta_{\xi\xi} = 0$$
 at $\xi = 0$ and 1



Adhesive Potential-Bubble Covered With Elastic Shell



 $\Delta \vec{F} = -\vec{\nabla}_{s} \cdot \left(\tau_{ss}\vec{t}_{s}\vec{t}_{s} + \tau_{\omega\omega}\vec{t}_{\omega}\vec{t}_{\omega} + q\vec{t}_{s}\vec{n}\right) + \vec{\nabla} \cdot \left[\left(W + \sigma\right)\vec{n}\vec{n}\right]$ $\vec{q} = \vec{\nabla}_s \cdot \underline{m} \cdot (\underline{I} - \vec{n}\vec{n})$ $W = W(y) = W_o \left| \left(\frac{\delta_A}{y} \right)^4 - 2 \left(\frac{\delta_A}{y} \right)^2 \right|$ Isothermal Gas Compression $P_{G0}V_0^{\gamma} = P_{G}V^{\gamma}$ Constitutive Law: Mooney-Rivlin $\tau_{ij} = \frac{G_s}{3\lambda_i\lambda_j} \left| \lambda_i^2 - \frac{1}{\lambda_i^2\lambda_i^2} \right| \left[1 + b(\lambda_j^2 - 1) \right]$ $m_i = k_b (K_i + vK_i) / \lambda_i, \quad K_i \equiv \lambda_i k_i - k_i^R$ *Effect of pre-stress* $\lambda_i = \frac{dS}{dS^{SF}}$, where $\lambda_i^0 = \frac{dS^0}{dS^{SF}} \neq 1$ **Boundary Conditions** $r_{\xi} = \vartheta_{\xi\xi} = 0$ at $\xi = 0$ and 1

Normal and Tangential Force Balance

 $\Delta \vec{F} = \Delta \vec{P} \Longrightarrow \Delta \vec{F}_{N} = (P_{G} - P_{atm}) I \cdot \vec{n} \quad \& \quad \Delta \vec{F}_{s} = \vec{0}$

[Blount et al., Proc. R. Soc. A, 2013]



Adhesive Potential-Free Bubble



$$P_{G} - P_{atm} = 2H(\sigma - W) + \frac{\partial W}{\partial n}$$
$$W = W(y) = W_{o} \left[\left(\frac{\delta_{A}}{y} \right)^{4} - 2 \left(\frac{\delta_{A}}{y} \right)^{2} \right]$$

Normal Force Balance

Arc-Length $ds^{2} = (rd\vartheta)^{2} + dr^{2}$

Isothermal Gas Compression $P_{G0}V_0^{\gamma} = P_G V^{\gamma}$

Boundary Conditions $r_{\xi} = \vartheta_{\xi\xi} = 0$ at $\xi = 0$ and 1

[Chamakos, Kavousakis & Papathanasiou, Soft Matter, 2013]



Finite Elements

Weak Form-Classic contact problem

$$\begin{aligned} R_{1} &= \int_{0}^{1} \left[\left(k_{s} \tau_{ss} + k_{\varphi} \tau_{\varphi\varphi} + \sigma 2H + P(\xi) + P_{atm} - P_{G} \right) B_{i} r \sin \vartheta s_{\xi} - \frac{\sigma \left(B_{i,\xi\xi} s_{\xi} - B_{i,\xi} s_{\xi\xi} \right)}{s_{\xi}^{2}} - \frac{m_{\varphi} B_{i,\xi} \sigma_{\xi}}{s_{\xi}} \right] d\xi + M_{BT} \\ R_{2} &= \int_{0}^{1} \left[\tau_{ss} B_{i,\xi} \sigma + B_{i} \sigma_{\xi} \tau_{\varphi\varphi} + \sigma m_{s} \left(k_{s,\xi} B_{i} + k_{s} B_{i,\xi} \right) + k_{s} m_{\varphi} B_{i} \sigma_{\xi} \right] d\xi + N_{BT} \\ R_{3} &= P_{G} V^{\nu} - P_{atm} V_{0}^{\nu} \\ \text{where} \\ \xi &= 0: \left(\vartheta = 0 \right) \quad M_{BT} = 0 \text{ and } N_{BT} = 0 \\ \xi &= 1: \left(\vartheta = \frac{\pi}{2} \right) M_{BT} = \frac{\sigma m_{s} B_{i,\xi}}{s_{\xi}} - \sigma q B_{i} \text{ and} \\ N_{BT} &= - \left[k_{s} m_{s} + \tau_{ss} \right] \sigma B_{i} \\ s_{\xi} &= \frac{ds}{d\xi} = \sqrt{r^{2} \vartheta_{\xi}^{2} + r_{\xi}^{2}} \quad ; \quad \sigma = r \sin \vartheta \quad ; \quad \sigma_{\xi} = \frac{d\sigma}{d\xi} = r_{\xi} \sin \vartheta + r \vartheta_{\xi} \cos \vartheta \end{aligned}$$



Finite Elements

Weak Form-Adhesive Potential with Elasticity

$$R_{1} = \int_{0}^{1} \left[\left(k_{s} \tau_{ss} + k_{\varphi} \tau_{\varphi\varphi} + (\sigma - W) 2H + P_{atm} - P_{G} \right) B_{i} r \sin \vartheta s_{\xi} - \frac{\sigma \left(B_{i,\xi\xi} s_{\xi} - B_{i,\xi} s_{\xi\xi} \right)}{s_{\xi}^{2}} - \frac{m_{\varphi} B_{i,\xi} \sigma_{\xi}}{s_{\xi}} + \frac{\partial W}{\partial n} \right] d\xi + M_{BT}$$

$$R_{2} = \int_{0}^{1} \left[\tau_{ss} B_{i,\xi} \sigma + B_{i} \sigma_{\xi} \tau_{\varphi\varphi} + \sigma m_{s} \left(k_{s,\xi} B_{i} + k_{s} B_{i,\xi} \right) + k_{s} m_{\varphi} B_{i} \sigma_{\xi} + \frac{\partial W}{\partial s} \right] d\xi + N_{BT}$$

$$R_{3} = P_{G} V^{\gamma} - P_{atm} V_{0}^{\gamma}$$

Weak Form-Adhesive Potential for a Free Bubble

$$R_{1} = \int_{0}^{1} \left[2H(\sigma - W) + P_{atm} - P_{G} + \frac{\partial W}{\partial n} \right] B_{i}r \sin \vartheta s_{\xi} d\xi$$
$$R_{2} = \int_{0}^{1} \left[S_{\max} - \left(r^{2}\vartheta_{\xi}^{2} + r_{\xi}^{2}\right)^{0.5} \right] B_{i}r \sin \vartheta s_{\xi} d\xi$$
$$R_{3} = P_{G}V^{\gamma} - P_{atm}V_{0}^{\gamma}$$



Finite Elements

- Method of Weighted Residuals: $\underline{\underline{J}} \cdot \vec{c} = \vec{R}$ $\underline{\underline{J}}$: Jacobian Matrix \vec{c} : Uknown coefficients \vec{R} : Residuals
- Basis Functions: B-Cubic Splines
- Calculation of Integrals:
- Solution of Linear System:
- Validation:

4 Points Gauss Quadrature

 $B_{i}(t_{j}) = \begin{cases} 4, & j = 1 \\ 1, & j = i \pm 1 \\ 0, & j = i \pm 2 \end{cases}$

Newton-Raphson Iterations Simple or Arc-Length Continuation

Calculation of the critical buckling load and comparison with previous analytical or experimental results.

> [Pelekasis, Manolis, Tsamopoulos, *Phys. Fluids,* 1990] [Tsiglkifis & Pelekasis, *Phys. Fluids,* 2011] [Lytra & Pelekasis, *Fluid Dyn. Res.,* 2014]



Continuation of the solution

Simple Continuation:

$$\lambda^{N+1} = \lambda^N + \Delta \lambda$$

Initial guess for the
new solution (N+1): $\vec{c}^{N+1} = \vec{c}^N$

 \vec{c}^{N+1} \vec{c}^N λN λ^{N+1} \vec{c}^N \vec{c}^{N+1} λ^N λ^{N+1} Limit point

Convergence

Arc-Length Continuation:

$$\begin{bmatrix} J & \vec{R}_{\lambda} \\ \vec{N}_{c} & N_{\lambda} \end{bmatrix} \begin{bmatrix} \vec{c} \\ \lambda \end{bmatrix} = \begin{bmatrix} \vec{R} \\ N \end{bmatrix} \text{ and } \begin{bmatrix} J & \vec{R}_{\lambda} \\ \vec{N}_{c} & N_{\lambda} \end{bmatrix} \begin{bmatrix} \partial \vec{c} / \partial s \\ \partial \lambda / \partial s \end{bmatrix} = \begin{bmatrix} \partial \vec{R} / \partial s \\ \partial N / \partial s \end{bmatrix}$$
$$N = \left(c_{i}^{N+1} - c_{i}^{N} \right)^{2} + \left(\lambda^{N+1} - \lambda^{N} \right)^{2} - \Delta s^{2}$$

Initial guess for the new solution (N+1):

$$\vec{c}^{N+1} = \vec{c}^N - \frac{\partial \vec{c}^N}{\partial s} \Delta s$$
$$\lambda^{N+1} = \lambda^N - \frac{\partial \lambda^N}{\partial s} \Delta s$$

Initial Guess -

Newton-Raphson ______ Iterations



Validation of Classic Contact Modeling





Classic Contact Modeling-AFM Data of Polymeric Bubbles



- Linear region flat shape.
- Non-linear region buckled-shell.
- The simulation can capture the experimental data.
- The amount of pre-stress explains the different response of experimental curves.
- The discrepancy in higher values of deformation is probably due to 3D-effects.

[Glynos et al., *Langmuir*, 2009] [Lytra, Pelekasis, Sboros, Glynos, Koutsos, *Procedia IUTAM*, 2015]

Asymptotic Analysis for Polymeric Bubbles

Point Load



Reissner's eq.

$$\underline{\mathsf{q.}} \quad F = \frac{4}{\left[3\left(1-v^2\right)\right]^{0.5}} \frac{Eh^2}{R_o} \Delta$$

[Reissner, J. Math. Phys., 1946]

Pogorelov's eq.
$$F = \left[\frac{3.56E^2h^5}{(1-v^2)R_o^2}\Delta\right]^{0.5}$$

[Pogorelov, Am. Math. Soc., 1988]

Plane Contact



$$\frac{\text{Stage I:}}{\left[3\left(1-v^{2}\right)\right]^{0.5}} \frac{Eh^{2}}{R_{o}} \Delta + 0.06543 \frac{Eh}{R_{o}} \Delta^{2}$$

$$\frac{\text{Stage II:}}{\left[3\left(1-v^{2}\right)\right]^{0.5}} \frac{Eh^{3}}{\left(1-v^{2}\right)R_{o}} \left(\frac{\Delta\sqrt{1-v^{2}}}{h}\right)^{0.5}$$

[Lytra, Pelekasis, Zafiropoulou, Zisis, Giannakopoulos, 8th GRACM Full Paper]



Elastic Properties Estimation - Polymeric Bubbles



The elastic properties can be estimated by a single measurement, the f-d curve.

[Glynos et al., *Langmuir*, 2009] [Lytra, Pelekasis, Sboros, Glynos, Koutsos, *Procedia IUTAM*, 2015]

Adhesive Potential-Bubble Covered With Lipid



[Buchner Santos et al., Langmuir, 2012]

Adhesive Potential-Bubble Covered With Lipid





- The shell remains flattened.
- The contact angle is not constant.



Asymptotic Analysis for Lipids





Adhesive Potential-Parametric Study (Elasticity Model)





- The classic contact model can easily simulate the force-deformation curve of a microbubble covered with polymeric shell.
 - The linear regime followed by a non-linear is associated with a sub-critical bifurcation.
 - In the linear regime, the shell remains flattened, but in the non-linear buckling takes place.
- In cases of lipids, which are softer, the intermolecular forces must be accounted for.
 - The experimental curve is almost linear, which indicate that buckling is not taking place.
 - Three models are considered in the present work, where no buckling was observed. The change of length is preferable, than bending, due to adhesion, which attaches the shell in cantilever.
 - The elasticity only model seems to be the more suitable for the simulations of the static response of CA microbubbles, covered with lipid.
 - The estimation of the elastic properties, requires further asymptotic analysis and perhaps different experimental set-up with side view photos.



Current & Future Work

- Simulation of a drop and a microbubble resting on a flat surface (Trapped microbubbles).
- Further asymptotic analysis and estimation of the elastic properties of bubbles covered with lipids.

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Thank you for your attention! Ευχαριστώ πολύ!

Questions?

Adhesive Potential-Parametric Study (Free Bubble)





Bifurcation Diagrams



[Lytra & Pelekasis, Fluid Dyn. Res., 2014]



Bifurcation Diagrams



 G_s =80 MPa k_b =3·10⁻¹⁴ Nm σ =0.051 N/m Mooney-Rivlin

[Lytra & Pelekasis, Fluid Dyn. Res., 2014]



- The static response of microbubbles, covered with an elastic shell, subject to a uniform pressure was investigated.
 - In this problem, the shell remains spherical for increasing the external overpressure, but its volume is decreasing.
 - When the overpressure reaches the critical load, the shell buckles into a symmetric or a non-symmetric shape. The buckling is indicated by one more negative eigenvalue. Following simple continuation, the post-buckling f-d curve is generated.
 - In every post-buckling curve a limit point was accounted and arc- length continuation was performed.
 - In the case of Gs=40 MPa, the first bifurcation point is dominated by the a symmetric eigenmode (P2), which leads to a oblate or prolate shapes.
 - In the case of Gs=80 MPa, the first bifurcation point is dominated by a nonsymmetric eigenmode (P3) and non-symmetric shapes.

